

# **Kuga-Satake map for arbitrary dimension**

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Workshop on holomorphic symplectic varieties

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## Hodge structures

**DEFINITION:** Let  $V_{\mathbb{R}}$  be a real vector space. **A (real) Hodge structure of weight  $w$**  on a vector space  $V_{\mathbb{C}} = V_{\mathbb{R}} \otimes_{\mathbb{R}} \mathbb{C}$  is a decomposition  $V_{\mathbb{C}} = \bigoplus_{p+q=w} V^{p,q}$ , satisfying  $\overline{V^{p,q}} = V^{q,p}$ . It is called **integer Hodge structure** if one fixes an integer lattice  $V_{\mathbb{Q}}$  or  $V_{\mathbb{Z}}$  such that  $V_{\mathbb{R}} = V_{\mathbb{Q}} \otimes_{\mathbb{Q}} \mathbb{R}$  or  $V_{\mathbb{R}} = V_{\mathbb{Z}} \otimes_{\mathbb{Z}} \mathbb{R}$ . A Hodge structure is equipped with  $U(1)$ -action, with  $u \in U(1)$  acting as  $u^{p-q}$  on  $V^{p,q}$ . **Morphism** of integer Hodge structures is a map which is  $U(1)$ -invariant and preserves the lattice.

**DEFINITION: Weak polarization** on a Hodge structure of weight  $w$  is a  $U(1)$ -invariant non-degenerate 2-form  $h \in V_{\mathbb{Q}}^* \otimes V_{\mathbb{Q}}^*$  (symmetric or antisymmetric depending on parity of  $w$ ). It is called **polarization** if it is in addition satisfies  $-(\sqrt{-1})^{p-q} h(x, \bar{x}) > 0$  for each non-zero  $x \in V^{p,q}$ .

**DEFINITION: Period space** of (weakly polarized) Hodge structure with dimensions  $\dim V^{p,q} = v_{pq}$  is the space of all decompositions  $V_{\mathbb{C}} = \bigoplus_{p+q=w} V^{p,q}$  such that the above conditions are satisfied.

**REMARK:** The period space, for odd weight, **is a complex manifold**. Indeed, the  $V^{q,p}$  spaces for  $q > p$  are determined by  $V^{p,q}$  uniquely, hence the period space is an open subspace in the space of  $k$ -tuples of subspaces  $V^{p,q} \subset V \otimes \mathbb{C}$ , with  $p < q$ . For even weight, **the period space for (weakly) polarized Hodge structures is again a complex manifold**.

## Hodge structures and homogeneous spaces

**EXAMPLE:** The **Hodge structure of weight 1** is a decomposition  $V_{\mathbb{C}} = V^{1,0} \oplus V^{0,1}$ , with  $\overline{V^{1,0}} = V^{0,1}$ . Clearly, the Hodge structures of weight 1 are in bijective correspondence with complex structures on  $V$ . Therefore, **the period space of Hodge structures of weight 1 is identified with  $GL(2n, \mathbb{R})/GL(n, \mathbb{C})$ .**

**EXAMPLE:** The **Hodge structure of K3 type** is a (weakly polarized) Hodge structure  $V_{\mathbb{C}} = \bigoplus_{\substack{p+q=2 \\ p,q \geq 0}} V^{p,q}$  of weight 2 with  $\dim V^{2,0} = 1$ .

**REMARK:** The **period space of weakly polarized Hodge structures of K3 type** is identified with the quadric of lines  $Q := \{l \in \mathbb{P}V_{\mathbb{C}} \mid h(l, l) = 0, h(l, \bar{l}) \neq 0\}$ .

### **THEOREM: (Kuga-Satake)**

Let  $Q$  be the space of weakly polarized Hodge structures of K3 type on  $(W, h)$ . Then **there exists a vector space  $V$  equipped with  $SO(W)$ -action and an  $SO(W)$ -equivariant embedding from  $Q$  to the space of Hodge structures of weight 1 on  $V$ .**

## Kuga-Satake embedding and Clifford modules

**THEOREM: (Kuga-Satake)** Let  $Q$  be the space of weakly polarized Hodge structures of K3 type on  $(W, h)$ . Then **there exists a vector space  $V$  equipped with  $SO(W)$ -action and an  $SO(W)$ -equivariant embedding from  $Q$  to the space of Hodge structures of weight 1 on  $V$ .**

**Proof. Step 1:** For any Hodge structure of K3 type, the corresponding action of  $\mathfrak{u}(1)$  is generated by a skew-symmetric matrix  $\mu$  of rank 2, acting trivially on the orthogonal complement to a 2-dimensional plane  $l = \langle \operatorname{Re} \Theta, \operatorname{im} \Theta \rangle$ , where  $\Theta$  is a generator of  $V^{2,0}$ , and acting as  $\begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$  on  $l$ .

**Step 2:** Let  $\mathcal{C}\ell(W)$  be the Clifford algebra of  $W$ , and  $V$  a space with  $\mathcal{C}\ell(W)$ -action (such space is called **Clifford module**). Using the standard embedding  $\mathfrak{so}(W) \subset \mathcal{C}\ell(W)$ , we can consider  $\mu$  as an element of  $\mathcal{C}\ell(W)$ . Then  $\mu^2 = -1$  in the Clifford algebra, and this gives a complex structure on  $V$ . ■

**REMARK:** Kuga and Satake were interested in **constructing an embedding of the symmetric spaces** associated with polarized Hodge structures of weight 1 and of K3 type.

## Hyperkähler manifolds

**DEFINITION:** (E. Calabi, 1978)

Let  $(M, g)$  be a Riemannian manifold equipped with three complex structure operators  $I, J, K : TM \rightarrow TM$ , satisfying the quaternionic relation

$$I^2 = J^2 = K^2 = IJK = -\text{Id}.$$

Suppose that  $I, J, K$  are Kähler. Then  $(M, I, J, K, g)$  is called **hyperkähler**.

**REMARK:** A hyperkähler manifold  $M$  is equipped with 3 symplectic forms  $\omega_I, \omega_J, \omega_K$ . The form  $\Omega := \omega_J + \sqrt{-1} \omega_K$  **is a holomorphic symplectic 2-form on  $(M, I)$** . ■

**THEOREM:** (Calabi-Yau) Let  $M$  be a compact, holomorphically symplectic Kähler manifold. Then  $M$  **admits a hyperkähler metric**, which is uniquely determined by the cohomology class of its Kähler form  $\omega_I$ .

*Hyperkähler geometry is essentially the same as holomorphic symplectic geometry*

## Kuga-Satake construction in arbitrary dimension

**REMARK:** Let  $M$  be a hyperkahler manifold Kuga-Satake construction **gives an embedding from  $H^2(M)$  to the second cohomology of a torus, compatible with the Hodge structure.** Indeed,  $W$  is embedded to  $\Lambda^2(V)$ , where  $V$  is a  $\mathcal{C}l(W)$ -module.

**THEOREM:** For any hyperkahler manifold  $M$  of complex dimension  $n$ , **there exists a torus  $T$  of dimension  $n + l$  and an embedding of cohomology space  $H^*(M) \mapsto H^{*+l}(T)$  which is compatible with the Hodge structures and the Poincare pairing.** Moreover, this embedding is compatible with an action of the Lie algebra generated by all Lefschetz  $sl(2)$ -triples on  $M$ .

**REMARK:** The corresponding map from the period space of  $M$  to the period space of  $T$  **coincides with the Kuga-Satake map.**

## Holomorphically symplectic manifolds

**DEFINITION:** A **holomorphically symplectic manifold** is a complex manifold equipped with non-degenerate, holomorphic  $(2,0)$ -form.

**REMARK:** Hyperkähler manifolds are holomorphically symplectic. Indeed,  $\Omega := \omega_J + \sqrt{-1} \omega_K$  is a holomorphic symplectic form on  $(M, I)$ .

**THEOREM:** (Calabi-Yau) A compact, Kähler, holomorphically symplectic manifold **admits a unique hyperkähler metric in any Kähler class.**

**DEFINITION:** For the rest of this talk, **a hyperkähler manifold is a compact, Kähler, holomorphically symplectic manifold.**

**DEFINITION:** A hyperkähler manifold  $M$  is called **of maximal holonomy**, or **simple**, or **IHS**, if  $\pi_1(M) = 0$ ,  $H^{2,0}(M) = \mathbb{C}$ .

**Bogomolov's decomposition:** Any hyperkähler manifold admits a finite covering which is a product of a torus and several maximal holonomy hyperkähler manifolds.

**Further on, all hyperkähler manifolds are assumed to be of maximal holonomy.**

## The Bogomolov-Beauville-Fujiki form

**THEOREM:** (Fujiki). Let  $\eta \in H^2(M)$ , and  $\dim M = 2n$ , where  $M$  is hyperkähler. **Then**  $\int_M \eta^{2n} = cq(\eta, \eta)^n$ , for some primitive integer quadratic form  $q$  on  $H^2(M, \mathbb{Z})$ , and  $c > 0$  an integer number.

**Definition:** This form is called **Bogomolov-Beauville-Fujiki form**. **It is defined by the Fujiki's relation uniquely, up to a sign.** The sign is determined from the following formula (Bogomolov, Beauville)

$$\lambda q(\eta, \eta) = \int_X \eta \wedge \eta \wedge \Omega^{n-1} \wedge \bar{\Omega}^{n-1} - \frac{n-1}{n} \left( \int_X \eta \wedge \Omega^{n-1} \wedge \bar{\Omega}^n \right) \left( \int_X \eta \wedge \Omega^n \wedge \bar{\Omega}^{n-1} \right)$$

where  $\Omega$  is the holomorphic symplectic form, and  $\lambda > 0$ .

**Remark:**  $q$  has signature  $(b_2 - 3, 3)$ . It is negative definite on primitive forms, and positive definite on  $\langle \Omega, \bar{\Omega}, \omega \rangle$ , where  $\omega$  is a Kähler form.

## Multi-dimensional BBF form

This is the original motivation for the present work.

1. BBF form is a weak polarization on  $H^2(M)$  compatible with all complex structures.
2. It is not hard to produce a weak polarization on  $H^k(M)$  compatible with all complex structures. However, there is no canonical choice.

**DEFINITION:** Let  $a, b \in H^{2k}(M)$ ,  $M$  Kähler of complex dimension  $2n$ , and  $q \in \text{Sym}^2(H^2(M)) \subset H^4(M)$  be the element corresponding to the BBF form. Then **the multi-dimensional BBF form** is  $a, b \longrightarrow \int_M a \wedge b \wedge q^{n-k}$ .

**CONJECTURE: It is non-degenerate.**

## Lie superalgebras

**DEFINITION:** A **graded vector space** is a space  $V^* = \bigoplus_{i \in \mathbb{Z}} V^i$ .

**REMARK:** If  $V^*$  is graded, the endomorphisms space  $\text{End}(V^*) = \bigoplus_{i \in \mathbb{Z}} \text{End}^i(V^*)$  is also graded, with  $\text{End}^i(V^*) = \bigoplus_{j \in \mathbb{Z}} \text{Hom}(V^j, V^{i+j})$

**DEFINITION:** An operator on a graded vector space is called **even (odd)** if it shifts the grading by even (odd) number. The **parity**  $\tilde{a}$  of an operator  $a$  is 0 if it is even, 1 if it is odd. We say that an operator is **pure** if it is even or odd.

**DEFINITION:** A **supercommutator** of pure operators on a graded vector space is defined by a formula  $\{a, b\} = ab - (-1)^{\tilde{a}\tilde{b}}ba$ .

**DEFINITION:** A **graded Lie algebra** (Lie superalgebra) is a graded vector space  $\mathfrak{g}^*$  equipped with a bilinear graded map  $\{\cdot, \cdot\} : \mathfrak{g}^* \times \mathfrak{g}^* \longrightarrow \mathfrak{g}^*$  which is graded anticommutative:  $\{a, b\} = -(-1)^{\tilde{a}\tilde{b}}\{b, a\}$  and satisfies **the super Jacobi identity**  $\{c, \{a, b\}\} = \{\{c, a\}, b\} + (-1)^{\tilde{a}\tilde{c}}\{a, \{c, b\}\}$

## Supersymmetry in Kähler geometry

Let  $(M, I, g)$  be a Kähler manifold,  $\omega$  its Kähler form. **On  $\Lambda^*(M)$ , the following operators are defined.**

0.  $d, d^*, \Delta$ , because it is Riemannian.
1.  $L(\alpha) := \omega \wedge \alpha$
2.  $\Lambda(\alpha) := *L*\alpha$ . It is easily seen that  $\Lambda = L^*$ .
3. The Weil operator  $W|_{\Lambda^{p,q}(M)} = \sqrt{-1} (p - q)$

**THEOREM:** These operators generate a Lie superalgebra  $\mathfrak{a}$  of dimension  $(5|4)$ , acting on  $\Lambda^*(M)$ . Moreover, the Laplacian  $\Delta$  is central in  $\mathfrak{a}$ , hence  $\mathfrak{a}$  also acts on the cohomology of  $M$ .

**REMARK:** This is a convenient way to summarize the Kähler relations and the Lefschetz'  $\mathfrak{sl}(2)$ -action.

## Supersymmetry in hyperkähler geometry

Let  $(M, I, J, K, g)$  be a hyperkaehler manifold,  $\omega_I, \omega_J, \omega_K$  its Kaehler forms. **On  $\Lambda^*(M)$ , the following operators are defined.**

0.  $d, d^*, \Delta$ , because it is Riemannian.

1.  $L_I(\alpha) := \omega_I \wedge \alpha$

2.  $\Lambda_I(\alpha) := *L_I * \alpha$ . It is easily seen that  $\Lambda_I = L_J^*$ .

3. Three Weil operators  $W_I|_{\Lambda^{p,q}(M,I)} = \sqrt{-1}(p-q)$ ,  $W_J|_{\Lambda^{p,q}(M,J)} = \sqrt{-1}(p-q)$ ,  $W_K|_{\Lambda^{p,q}(M,K)} = \sqrt{-1}(p-q)$

**THEOREM:** These operators generate a Lie superalgebra  $\mathfrak{a}$  of dimension  $(11|8)$ , acting on  $\Lambda^*(M)$ . Moreover, the Laplacian  $\Delta$  is central in  $\mathfrak{a}$ , hence  $\mathfrak{a}$  also acts on the cohomology of  $M$ .

**REMARK:** The Weil operators form the Lie algebra  $\mathfrak{su}(2)$  of unitary quaternions. This means that **the quaternionic action belongs to  $\mathfrak{a}$** . In particular,  $L_J, L_K, \Lambda_J$  and  $\Lambda_K$ .

**REMARK:** The twisted de Rham differentials  $d_I, d_J, d_K$ , associated to  $I, J, K$  also belong to  $\mathfrak{a}$ :  $d_I = [W_I, d]$ ,  $d_J = [W_J, d]$ ,  $d_K = [W_K, d]$

**$\mathfrak{so}(4, 1)$ -action and the Hodge decomposition**

**REMARK:** 1.  $[L_I, \Lambda_J] = W_K$ ,  $[L_J, \Lambda_K] = W_I$ ,  $[L_I, \Lambda_K] = -W_J$ .

2. The even part of  $\mathfrak{a}$  **is isomorphic to**  $\mathfrak{sp}(1, 1, \mathbb{H}) \oplus \mathbb{R} \cdot \Delta$ .

3. The odd part  $\langle d, d_I, d_J, d_K, d, {}^*d_I^*, d_J^*, d_K^* \rangle$  **generates the 9-dimensional odd Heisenberg algebra**, with the only non-trivial supercommutators being  $\{d, d^*\} = \{d_I, d_I^*\} = \{d_J, d_J^*\} = \{d_K, d_K^*\} = \Delta$

4. The action of  $\mathfrak{a}_{\text{even}}$  on  $\mathfrak{a}_{\text{odd}}$  **is the fundamental representation of**  $\mathfrak{sp}(1, 1, \mathbb{H})$  **in**  $\mathbb{H}^2$ , with the quaternionic Hermitian metric on  $\mathfrak{a}_{\text{odd}}$  provided by the anticommutator.

**COROLLARY:** The weight decomposition of the  $\mathfrak{sp}(1, 1, \mathbb{H}) = \mathfrak{so}(4, 1)$ -action on  $H^*(M)$  **coincides with the Hodge decomposition.**

## Lefschetz-Frobenius algebras

**DEFINITION: A Frobenius algebra** is a graded commutative algebra  $A = \bigoplus_{i=0}^d A^i$  equipped with the Poincare-type non-degenerate product.

**DEFINITION:** A **Lefschetz triple** in a Frobenius algebra  $A = \bigoplus_{i=0}^{2n} A^i$  is a triple of operators  $L_\eta, H, \Lambda_\eta$  where  $\eta \in A^2$  is a fixed element,  $L_\eta(x) := \eta \wedge x$ ,  $H|_{A^i} = i - n$  and  $\Lambda_\eta$  is an element such that  $L_\eta, H, \Lambda_\eta$  is an  $\mathfrak{sl}(2)$ -triple. A Frobenius algebra admitting a Lefschetz triple is called **a Lefschetz-Frobenius algebra** (Looijenga, Lunts).

**REMARK:** Such  $\Lambda_\eta$  **is uniquely determined by  $H$  and  $\eta$**  (this statement is sometimes called “Morozov’s lemma”, and sometimes included in the statement of Jacobson-Morozov theorem).

**REMARK:** Existence of  $\Lambda_\eta$  for given  $\eta \in A^2$  is an open property in  $A^2$ , hence **a Lefschetz-Frobenius algebra admits many  $\mathfrak{sl}(2)$ -triples.**

## Lia algebra $\mathfrak{g}$ generated by $\mathfrak{sl}(2)$ -triples

**THEOREM:** Let  $M$  be a hyperkähler manifold of maximal holonomy,  $A^*$  its cohomology algebra and  $\mathfrak{g} := \mathfrak{g}(A)$  the Lie algebra generated by all Lefschetz  $\mathfrak{sl}(2)$ -triples. **Then  $\mathfrak{g}$  is isomorphic to  $\mathfrak{so}(b_2 - 2, 4)$ .**

**Sketch of the proof. Step 1:** Consider the action of  $\mathfrak{g}$  on the **Mukai extension**  $\hat{H}^2(M) := \mathbb{R} \cdot x \oplus H^2(M) \oplus \mathbb{R} \cdot y$ , where  $x$  has grading 0,  $y$  has grading 4,  $H^2(M)$  has grading 2. We equip  $\hat{H}^2(M)$  with **the Mukai form** which is equal to BBF on  $H^2(M)$ , preserves grading, and satisfies  $q_M(x, y) = 1$ ,  $x^2 = y^2 = 0$ ,  $x, y \perp H^2(M)$  and  $(x, y) = 1$ . The action of  $\mathfrak{g}$  on  $\hat{H}^2(M)$  is determined by the following properties: **1. It is compatible with the grading. 2. For all  $\alpha, \beta \in H^2(M)$ , one has  $L_\alpha x = \alpha$ ,  $L_\alpha \beta = q(\alpha, \beta)y$ , where  $q$  is the BBF form. 3.  $\Lambda_\alpha y = \alpha$ ,  $\Lambda_\alpha \beta = q(\alpha, \beta)x$ .**

To see that this action is well-defined, we need to check that commutator relations hold. This follows from commutator relations in  $\mathfrak{so}(1, 4)$  and Zariski density of pairs  $\alpha, \beta \in \langle \omega_I, \omega_J, \omega_K \rangle$  in the set of all pairs  $\alpha, \beta \in H^2(M)$ .

**Step 2:** The map  $\mathfrak{g} \rightarrow \mathfrak{so}(\hat{H}^2(M))$  is surjective, which follows from the dimension argument (dimensions are computed using the local Torelli theorem). Injectivity of  $\mathfrak{g} \rightarrow \mathfrak{so}(\hat{H}^2(M))$  is clear, because  $\mathfrak{so}(\hat{H}^2(M))$  is given by generators and relations which hold true in  $\mathfrak{g}$ . ■

## Hodge structures and $\mathfrak{g}$ -action

**REMARK:** The Lie algebra  $\mathfrak{g} = \mathfrak{so}(b_1 - 2, 4)$  is equipped with a grading  $\mathfrak{g} = \mathfrak{g}_{-2} \oplus \mathfrak{g}_0 \oplus \mathfrak{g}_2$ , induced by the grading on the Mukai space:  $\hat{H}^2(M) := H_0 \oplus H^2(M) \oplus H_4$ , with  $H_0$  and  $H_4$  1-dimensional. Then  $\mathfrak{g}_0 = \mathfrak{g}'_0 \oplus H$ , where  $H = [L_\omega, \Lambda_\omega]$  is the operator inducing the grading and commuting with the rest of  $\mathfrak{g}_0$ , denoted by  $\check{\mathfrak{g}}_0$ .

**REMARK:** The Lie algebra  $\mathfrak{g}'_0 := \mathfrak{so}(b_1 - 1, 3)$  is generated by the Weil maps  $W_I$  for all complex structures  $I$  of hyperkähler type obtained by deformations. The corresponding Lie group  $G_0$  acts as  $\text{Spin}(b_1 - 1, 3)$  in odd-dimensional cohomology and  $\text{SO}(b_1 - 1, 3)$  on even-dimensional ones. It is generated by the complex structure action on  $H^2(M)$  for all deformations of  $I$ .

**COROLLARY:** Let  $M$  be a hyperkähler manifold, and  $H^*(M) \hookrightarrow H^{*+l}(T)$  an embedding to the cohomology of a torus. Suppose that this embedding is compatible with an action of the Lie algebra generated by all Lefschetz  $\mathfrak{sl}(2)$ -triples on  $M$ . **Then it is compatible with the Hodge structures**, in the same sense as the usual Kuga-Satake map.

## $k$ -symplectic structures

**DEFINITION:** Let  $V$  be a  $4n$ -dimensional vector space, and  $\Psi : W \rightarrow \Lambda^2(V)$  a linear map. Assume that  $\Psi(\omega)$  is a symplectic form for general  $\omega \in W$ , and has rank  $\frac{1}{2} \dim W$  for  $\omega$  in a non-degenerate quadric  $Q \subset W$ . Then  $\Psi$  is called  **$k$ -symplectic structure on  $V$** , where  $k = \dim W$ .

**EXAMPLE:** For  $k = 3$  this is **hypersymplectic structure**, which is the same as a triple of symplectic form  $\omega_1, \omega_2, \omega_3$  such that  $\omega_i \omega_j^{-1}$  act on  $V$  as the matrices algebra  $\text{Mat}(2)$ . In a sense, hypersymplectic structure is “a complexification of a hyperkähler structure”.

**THEOREM: (Soldatenkov-V.)** Let  $V$  be a  $k$ -symplectic space. **Then  $V$  is a Clifford module over a Clifford algebra  $\mathcal{Cl}(W)$ .**

**THEOREM:** In this situation, **the algebra generated by  $\mathfrak{sl}(2)$ -triples associated with  $\omega \in W$  acts on  $\Lambda^*(V)$  as the algebra  $\mathfrak{g} = \mathfrak{so}(\hat{W})$** , where  $\hat{W} = W \oplus \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$  is Mukai extension of  $W$ .

**Proof:** We need to check that the  $\mathfrak{so}(4, 1)$ -relations hold for any non-degenerate 3-dimensional subspace  $W_3 \subset W$ . However, this subspace is associated with a hypersymplectic structure, which is a complexification of a hyperkähler, and satisfies the same relations. ■

## Proof of the main result

**THEOREM:** For any hyperkahler manifold  $M$  of complex dimension  $n$ , **there exists a torus  $T$  of dimension  $n + k$  and an embedding of cohomology space  $H^*(M) \mapsto H^{*+l}(T)$  which is compatible with the Hodge structures and the Poincare pairing.** Moreover, this embedding is compatible with an action of the Lie algebra generated by all Lefschetz  $\mathfrak{sl}(2)$ -triples on  $M$ .

**Proof:** Let  $\mathfrak{g}$  be the Lie algebra generated by all  $\mathfrak{sl}(2)$ -triples, and  $W := H^2(M)$ . For any Clifford module  $V$  over  $\mathcal{Cl}(W)$ ,  $V$  admits a  $b_2$ -symplectic structure which gives  $\mathfrak{g}$ -action in  $\Lambda^*(V)$ . **If we manage to produce an embedding of  $\mathfrak{g}$ -modules  $H^*(M) \hookrightarrow \Lambda^*(V)$ , we are done.**

However,  $\Lambda^*(V)$  is an exact representation of  $\text{Spin}(\widehat{W})$ , hence its tensor powers contain any representation of  $\text{Spin}(\widehat{W})$ . These tensor powers correspond to  $\Lambda^*(V^n)$ , which is also a Grassmann algebra for a Clifford module over  $\mathcal{Cl}(W)$ .

■